

# Possibilities for a Global-Flagship **Ultra-Cold (and Very-Cold) Neutron Source at the SNS**

Kent Leung

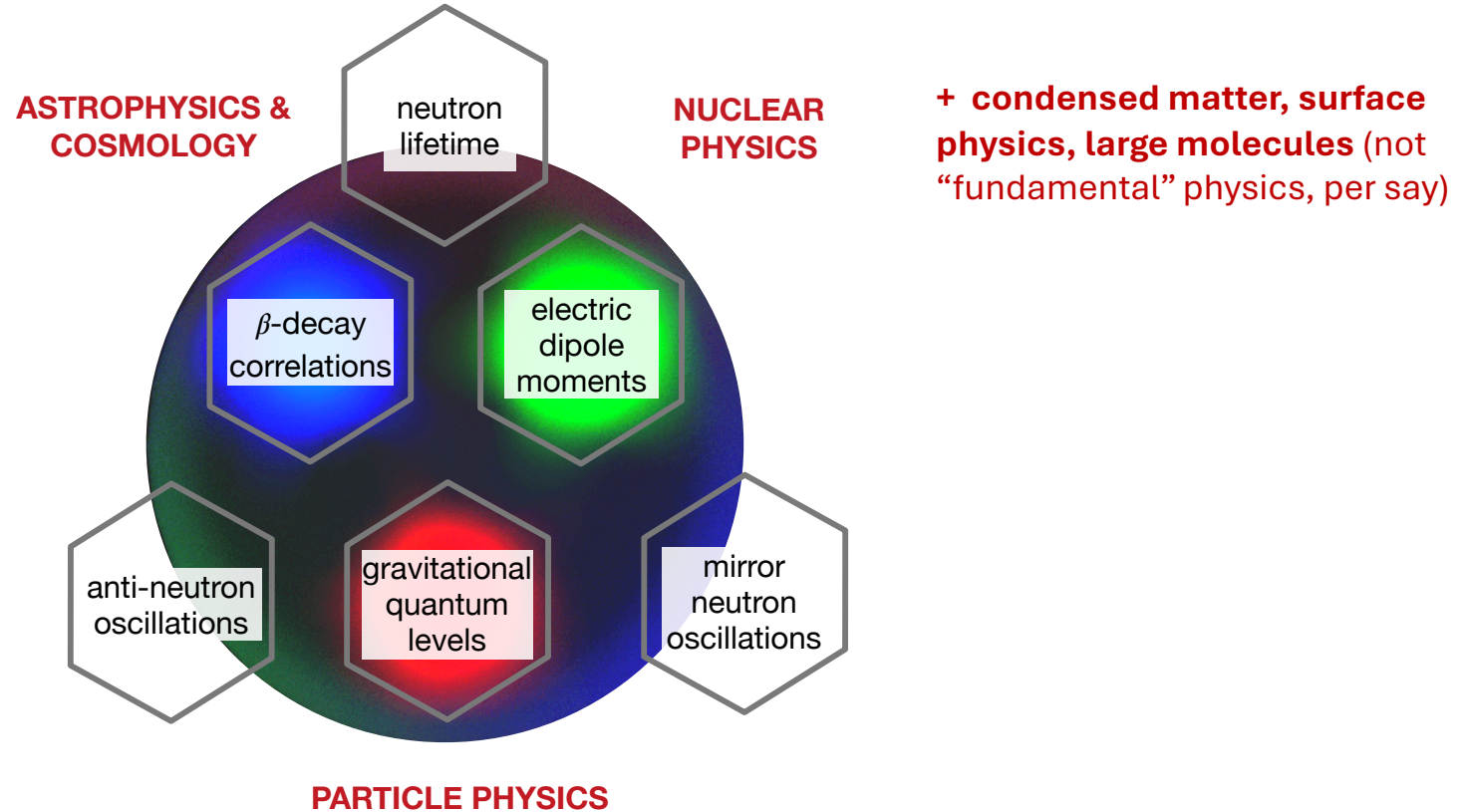
Physics & Astronomy Department, Montclair State University, NJ (tri-state area)



# Fundamental Neutron Physics

- Experimenter chooses between **all four fundamental forces** to study/control neutrons
- “Fundamental” **spin-1/2** baryon; **building block** of nuclei; **sizable magnetic moment**; unaffected by Lorentz force

## Common experiments:



Dubbers and Schmidt. The neutron and its role in cosmology and particle physics. *Rev. Mod. Phys.*, 83:1111 (2011)

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Common experiments:

**Ultracold neutron (UCN) experiments:**

*UCNs leads*

*UCNs competitive*

*UCNs only*

**ASTROPHYSICS & COSMOLOGY**

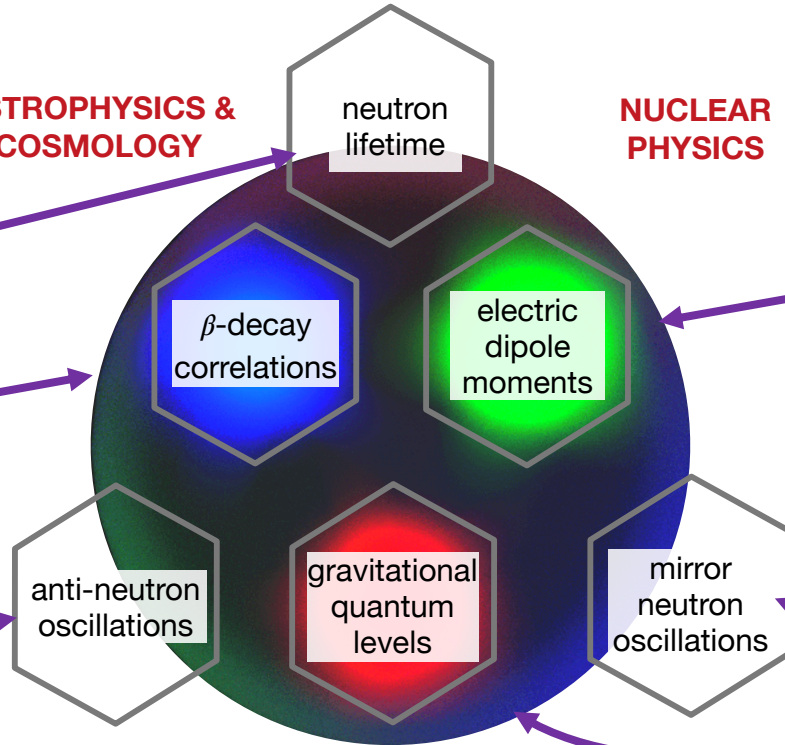
**NUCLEAR PHYSICS**

**+ condensed matter, surface physics, large molecules** (not “fundamental” physics, per say)

*UCN significantly leads*

*UCN possibilities*  
(Some experiments already + new theoretical suggestions)

**PARTICLE PHYSICS**



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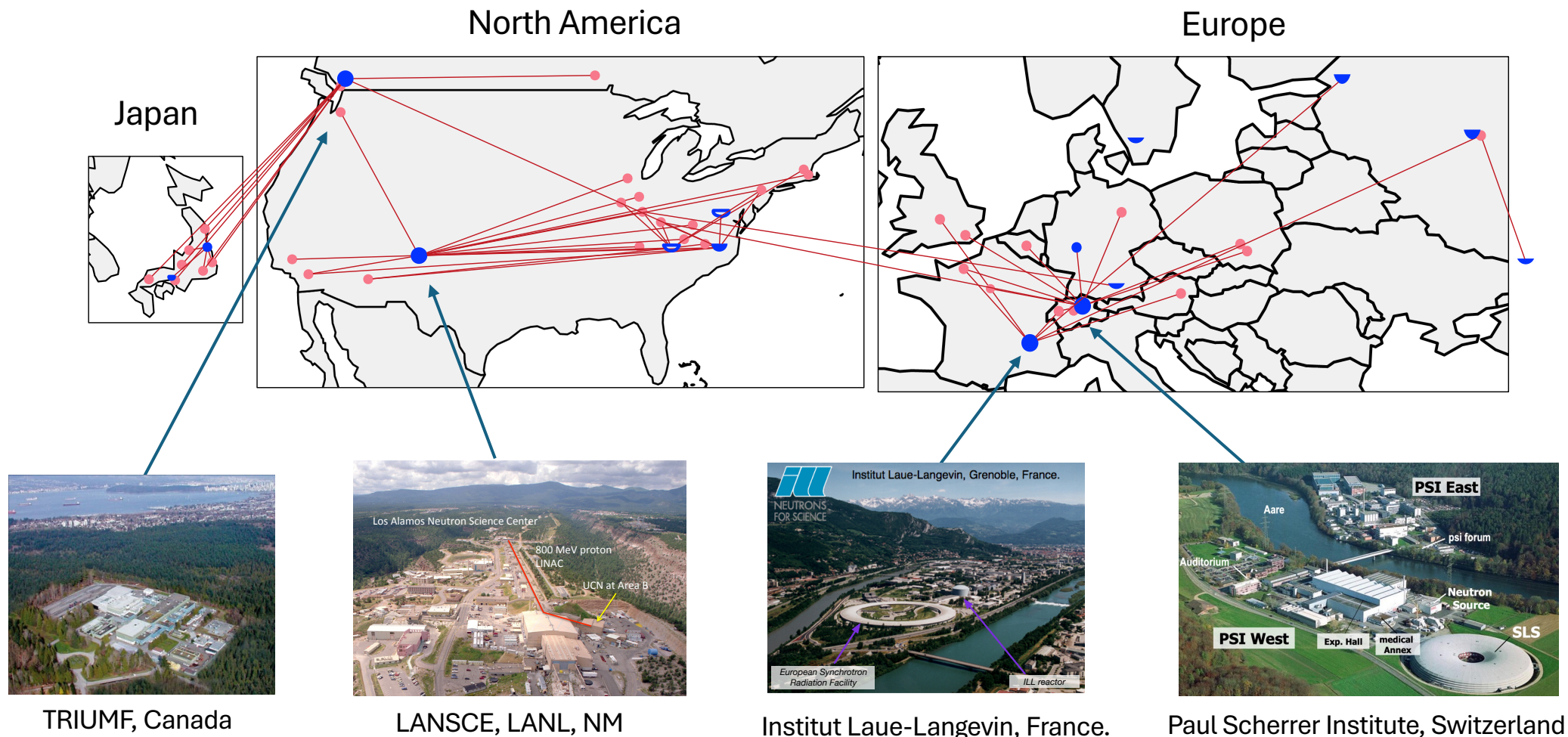
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# Ultracold Neutron Experimental Landscape

● = major operating sources: LANL, TRIUMF, ILL and PSI

● = collaborators

● = smaller operating sources: J-PARC, U. Mainz, ◐ = previous/proposed in-situ UCN experiments, ◑ = proposed/under construction



TRIUMF, Canada

LANSCE, LANL, NM

Institut Laue-Langevin, France.

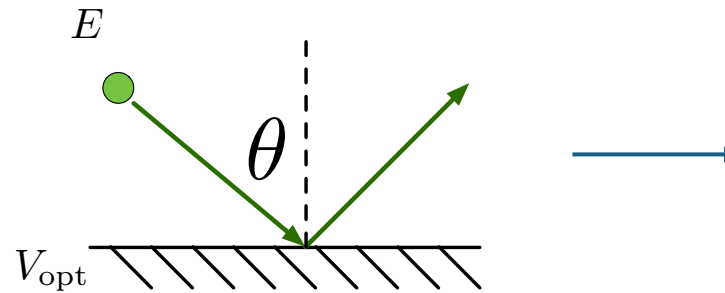
Paul Scherrer Institute, Switzerland

# What are ultracold neutrons?

Neutron optical or  
“Fermi” potential

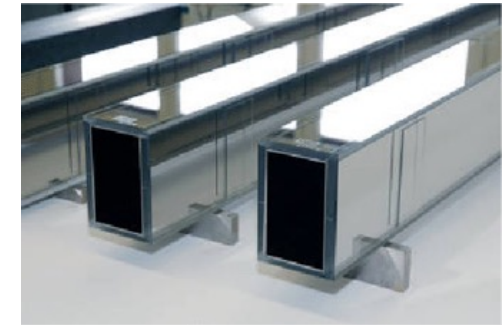
Material	$V_{\text{opt}}$
$^{58}\text{Ni}$	335 neV
Be	252 neV
Fluorocarbons	~100 neV
Al	54 neV
polyethylene	-9 neV
d8-polystyrene	160 neV

Reflection occurs if:  $E_{\perp} = E \cos^2 \theta \leq V_{\text{opt}}$



(Reflection can be non-specular/diffuse)

Neutron optics

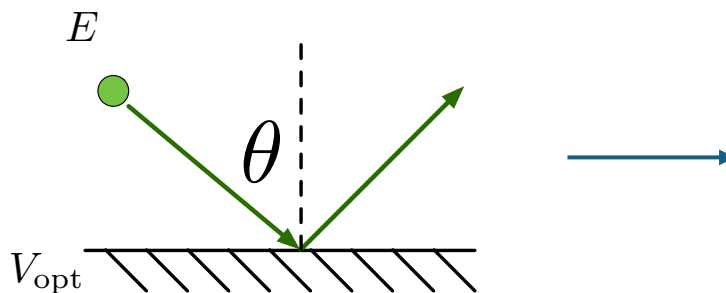


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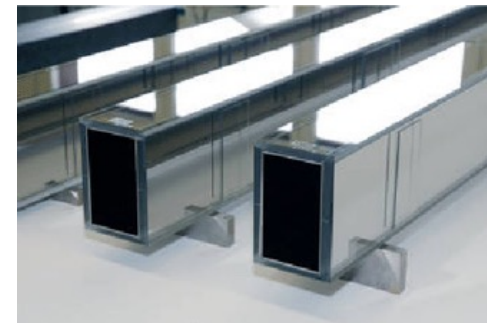
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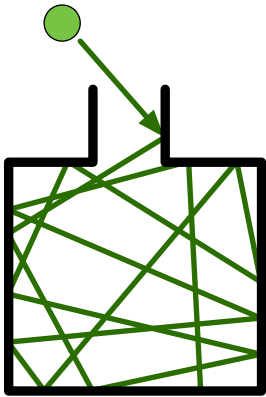
**If energy sufficiently low (“ultracold”)**  
**→ optical reflections at all incident angles**  
**→ ultracold neutrons can be stored in a box**

*First idea: Fermi (1946); Published: Zeldovich (1959).*

# UCN energy, treatment and losses

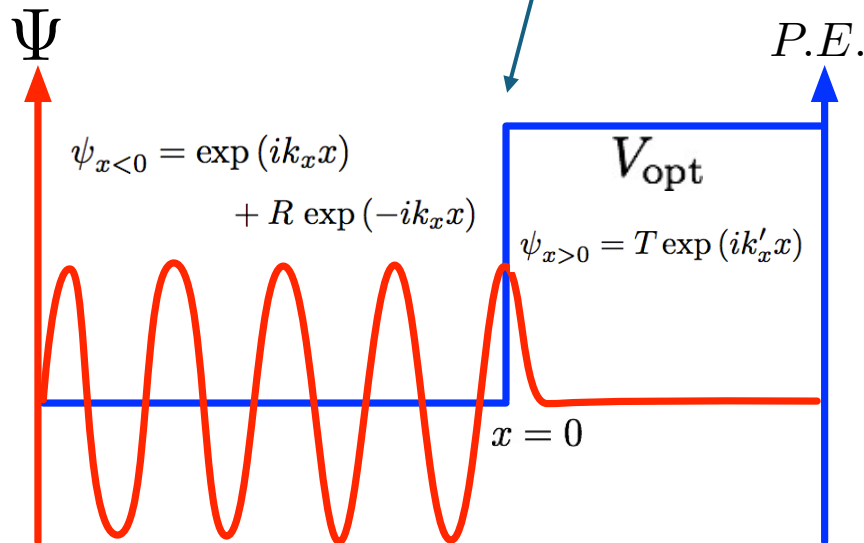
**UCNs:  $E < 300$  neV ,  $\lambda > 60$  nm,  $v < 7$  m/s, “temperature”  $< 2$  mK**

- “ideal gas” of neutrons but not in thermal equilibrium with material (typically at 300 K)
- Treat UCN reflection as 1D particle incident on a potential barrier
- UCN wavefunction tunnels into barrier so some loss during reflections. Losses treated as imaginary part of optical potential:

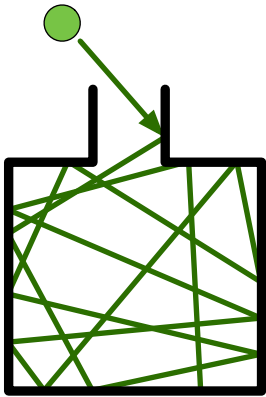


$$U_{\text{opt}} = V_{\text{opt}} + iW$$

$$W = \frac{\hbar}{2} \sum_i n_i \sigma_{\text{loss}}(v) v$$



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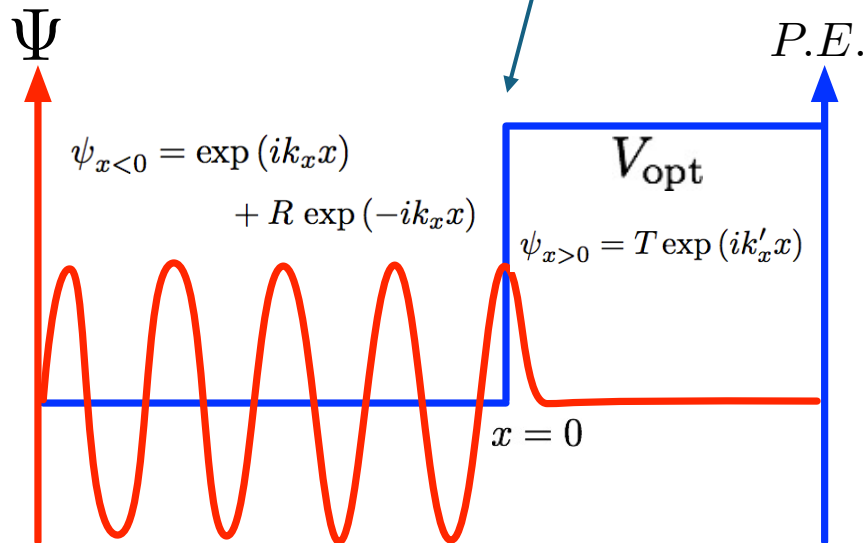


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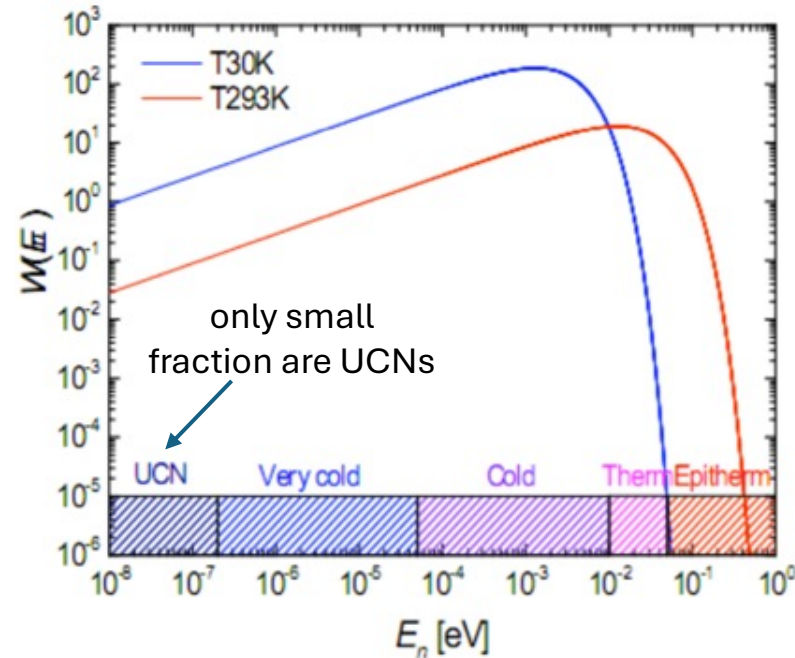


## Loss mechanisms

- nuclear capture
- **up-scattering off excitations** (e.g. phonons) in material to above UCN energy
- **Typically  $10^{-4}$  loss probability per reflection**  $\Rightarrow$  store UCNs with time constant  $\tau_{\text{tot}} \lesssim 600$  s depending on size of bottle
- Can use loss rate to understand detailed properties of surfaces also (surface/condensed matter)

# Producing UCNs: the limits of moderation

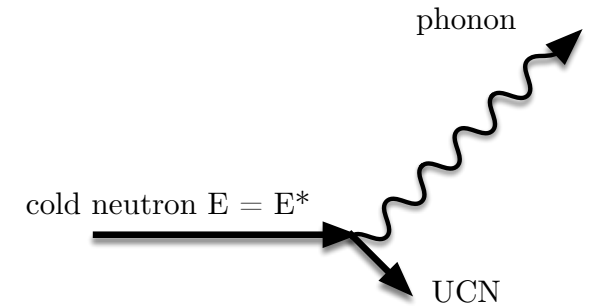
- **Fast neutrons (~MeV)** moderated in a material by scattering off protons and deuterons
- Neutrons come to **thermal equilibrium** in material & follow **Maxwell-Boltzmann distribution**:



- With moderation, UCN density  $\rho_{\text{UCN}}$  scales as  $\sim T^{-3/2}$
- The estimated  $\rho_{\text{UCN}}$  inside LD<sub>2</sub> cold moderators at most powerful neutron sources  $\sim 1000$  UCN/cm<sup>3</sup>
- Can't increase density further with conservative potentials due to **Liouville's theorem on phase space density**. And then there are **UCN transport losses to an experiment**.
- ILL's PF2 extracts out of cold moderator to experiment achieves  $\sim 50$  UCN/cm<sup>3</sup>

# “Super-thermal” production of UCNs

- **Scattering off phonons to convert cold neutrons to UCN energy range.**  
Phonons pumped away by refrigerator. [Golub & Pendlebury, Physics Letters A (1977)]
- Two competing materials: **superfluid  $^4\text{He}$**  (He-II) and **solid deuterium ( $\text{SD}_2$ )**



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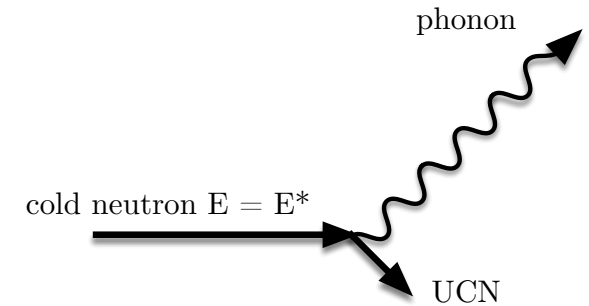
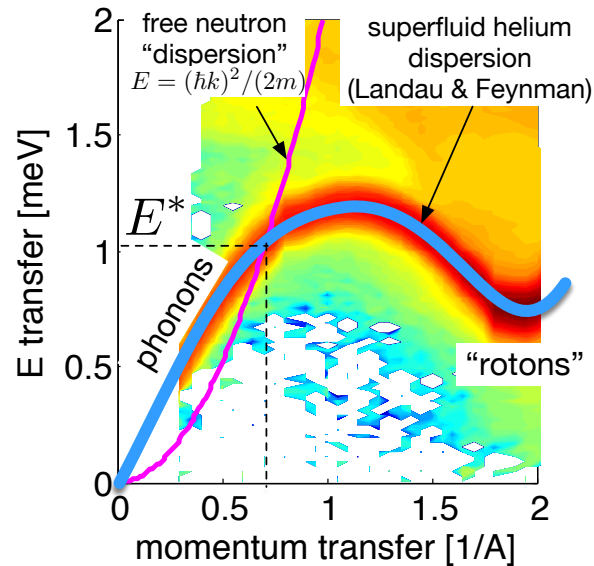
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- Consider He-II:

Log contour plot of dynamic structure factor of He-II @ 1.2K  
from [Andersen et al. J. Phys. Condens. Matter (1994)]



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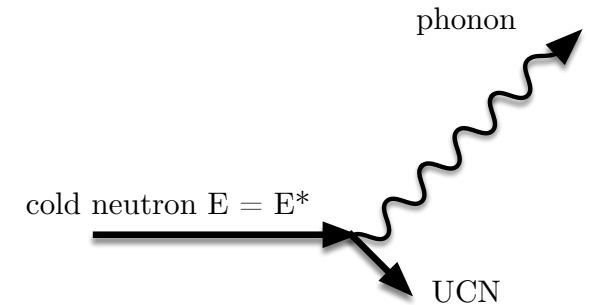
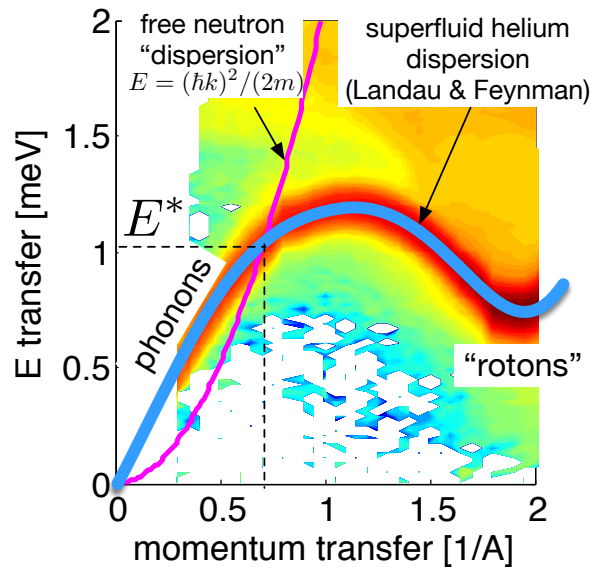
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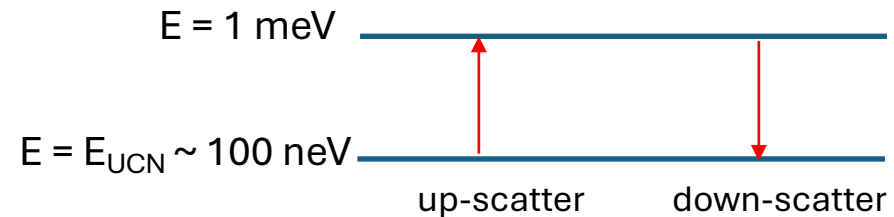
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1 meV ( $9\text{\AA}$ ) CN & UCN state form a two level system:



**In equilibrium the two processes linked by detailed balance:**

$$\sigma(E_{UCN} \rightarrow E_{UCN} + E^*) = \frac{E_{UCN} + E^*}{E_{UCN}} e^{-E^*/k_B T} \sigma(E_{UCN} + E^* \rightarrow E_{UCN})$$

up-scattering
suppressed by Boltzmann factor (i.e. exponential)
down-scattering ( $\sim$  constant)

- For  $T = 0.5\text{K}$ , **up-scattering time constant  $\approx 2$  hours.** Neutron absorption by  $^4\text{He}$  is zero.
- “Super-thermal” because UCNs are not in equilibrium with He-II converter

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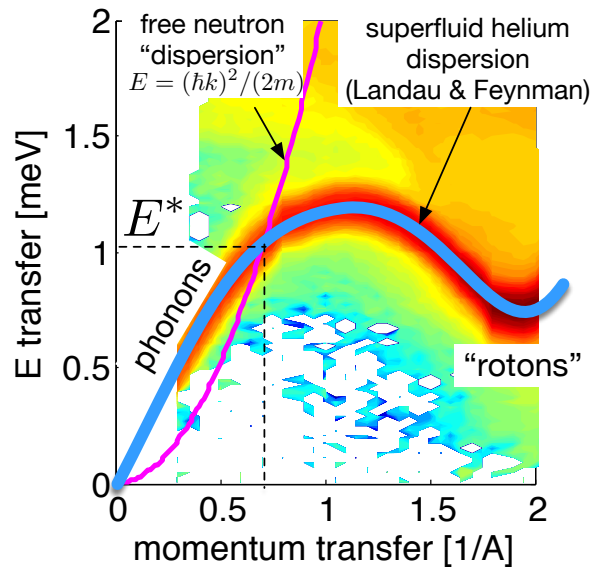
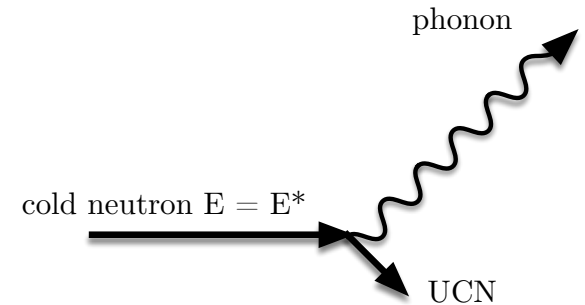
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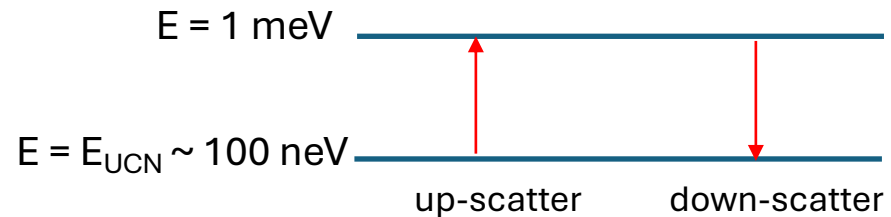
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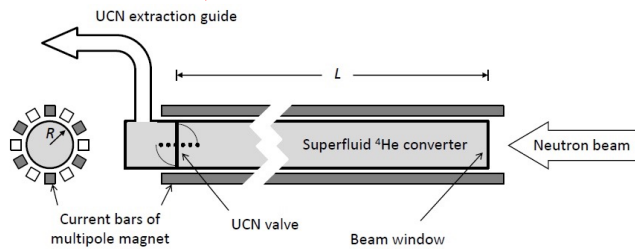
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- “Super-thermal” because UCNs are not in equilibrium with He-II converter
- **$\text{SD}_2$  has larger production cross-section (many phonon modes), only requires cooling to  $\sim$  6 K, but up-scattering time  $\sim$ 50 ms**

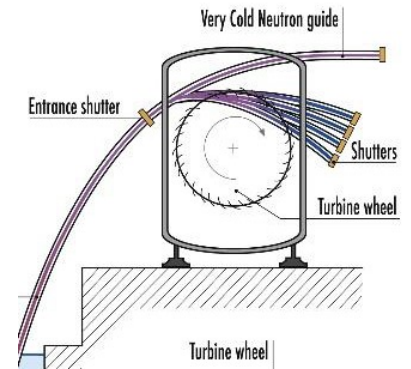
# UCN density or current source

Density [UCN/cm<sup>3</sup>]

Current [UCN/s]  
(or “integrated flux”)



ILL's PF2 Doppler turbine (~10 UCN/cm<sup>3</sup> @ experiment).



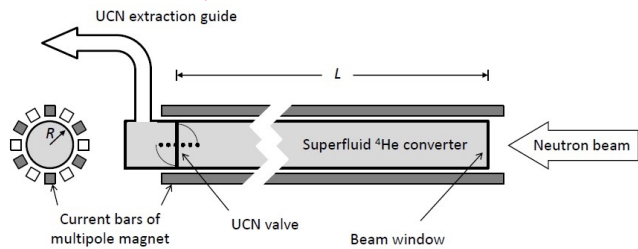
*PF2 is also a great VCN source...*

ILL SUN-series: ~0.5 K He-II  
at end of beam (**aim ~100  
UCN/cm<sup>3</sup>** @ nEDM  
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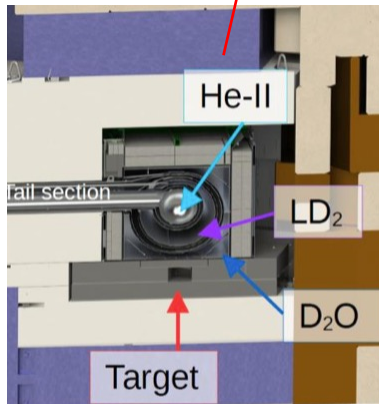
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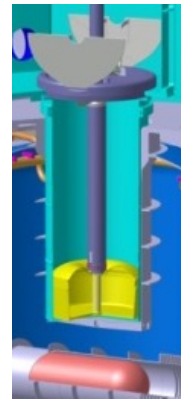


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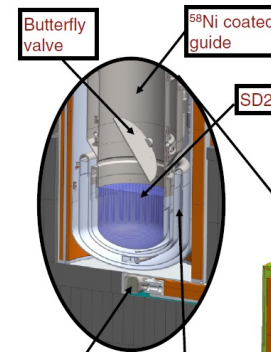


TRIUMF in-pile ~1 K  
4He. **Projected**  
based on 2025  
results: **~ 700  
UCN/cm<sup>3</sup>** @ nEDM  
experiment

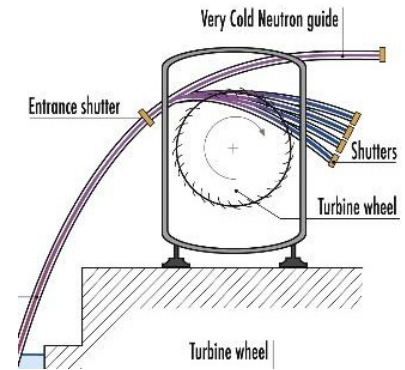
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LANL sD2 (~ 100  
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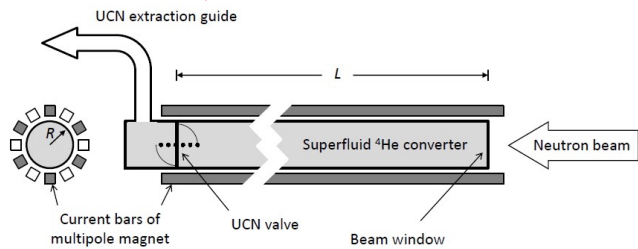


PNPI in-pile 4He (~1.3 K)

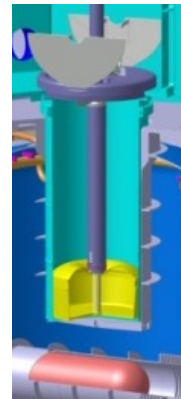
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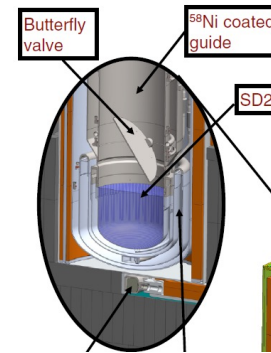
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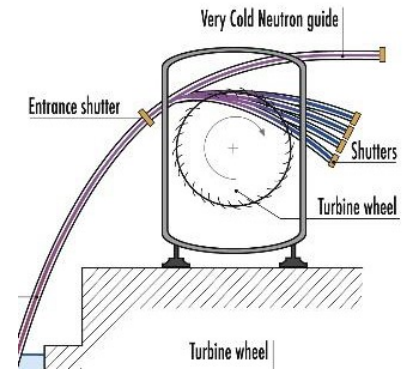
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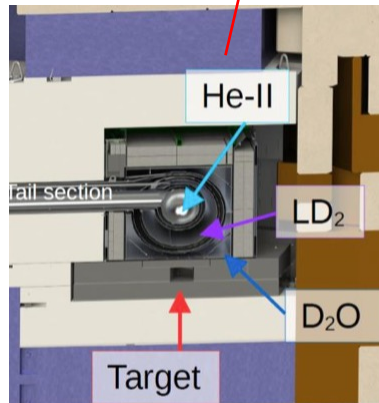


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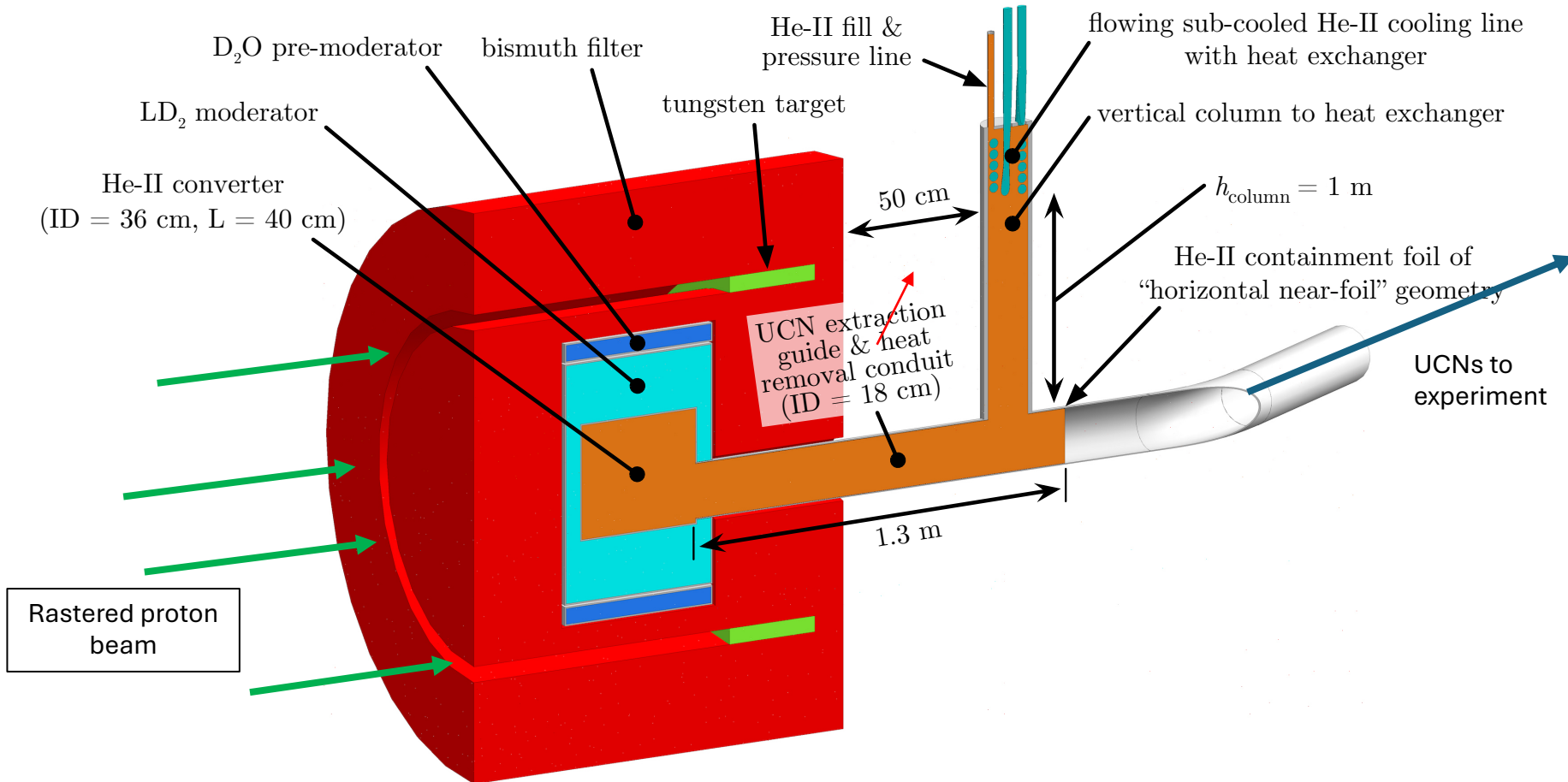
PNPI in-pile 4He (~1.3 K)

What about a **He-II  
current-optimized  
UCN source?**!

# A next-generation inverse-geometry spallation-driven UCN source

J. Appl. Phys. **126**, 224901 (2019); <https://doi.org/10.1063/1.5109879>

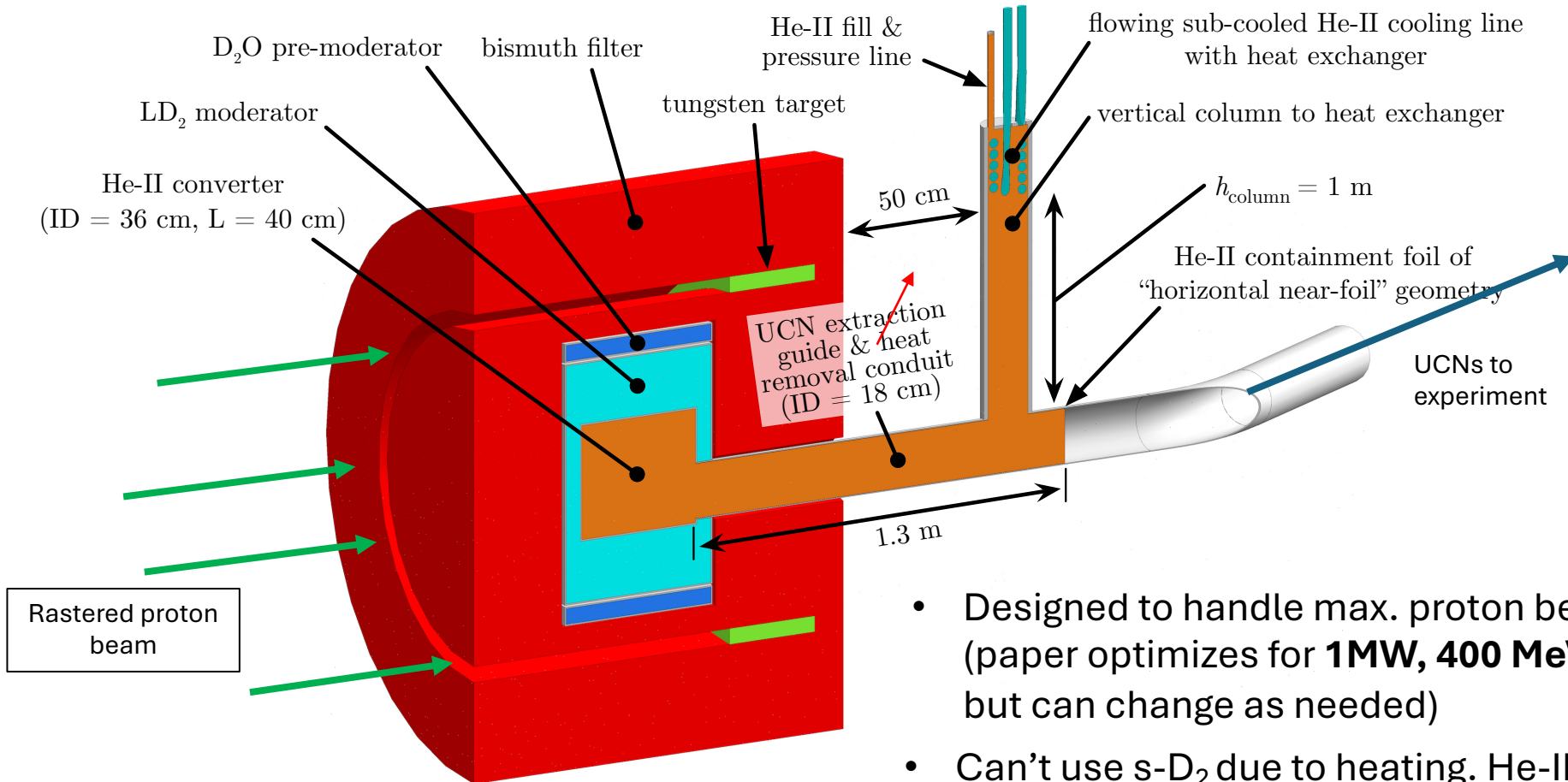
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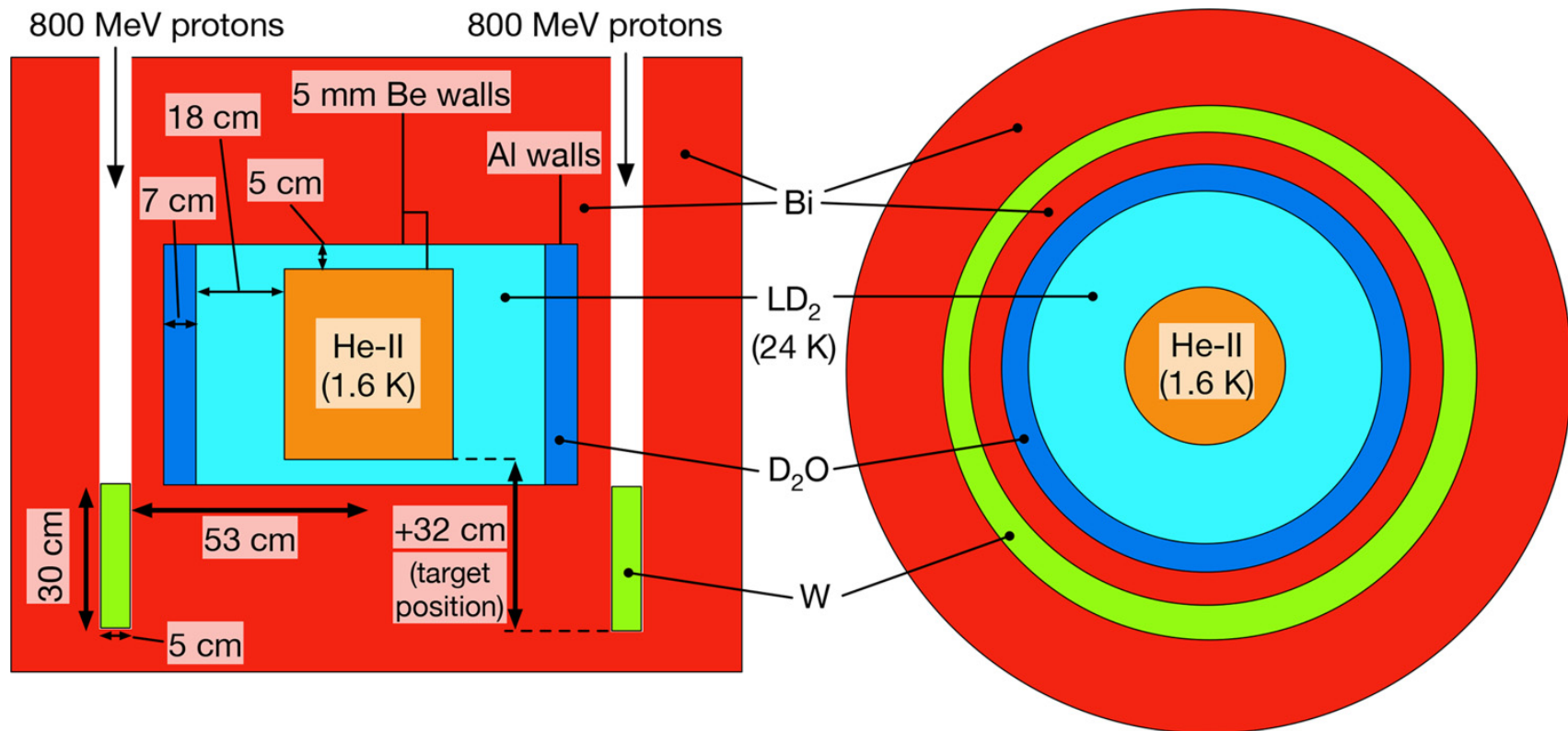
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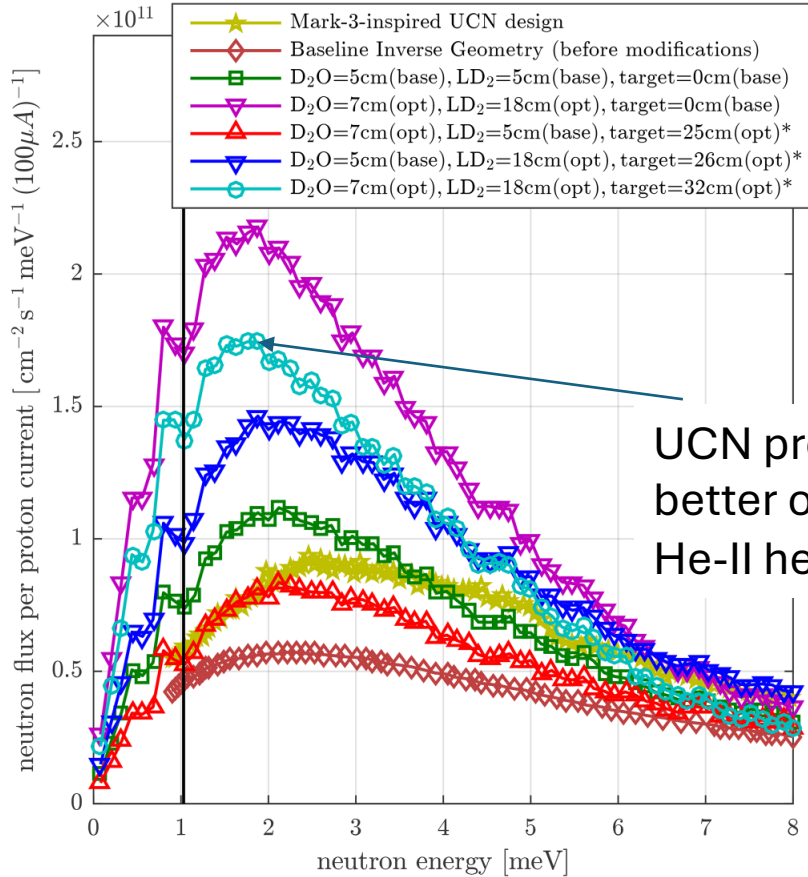
- Designed to handle max. proton beam powers (paper optimizes for **1MW, 400 MeV protons** but can change as needed)
- Can't use s-D<sub>2</sub> due to heating. He-II is a (near) perfect thermal conductor
- Leverage "recent" (~ 20 years) "**off-the-shelf**" **sub-cooled helium technology**

# Maximizing flux of 8.9Å / 1 meV neutrons

- Raster proton beam with kicker magnets to distribute heat in tungsten target => water cooling
- MCNP calculations bench-marked with Los Alamos' Lujan Center Mark-3 target
- Place **40-L volume of superfluid  $^4\text{He}$**  inside the target, pre-moderator, and moderator in “inverse geometry”
- Optimized D2O pre-moderator thickness, LD<sub>2</sub> moderator thickness, and “target location”

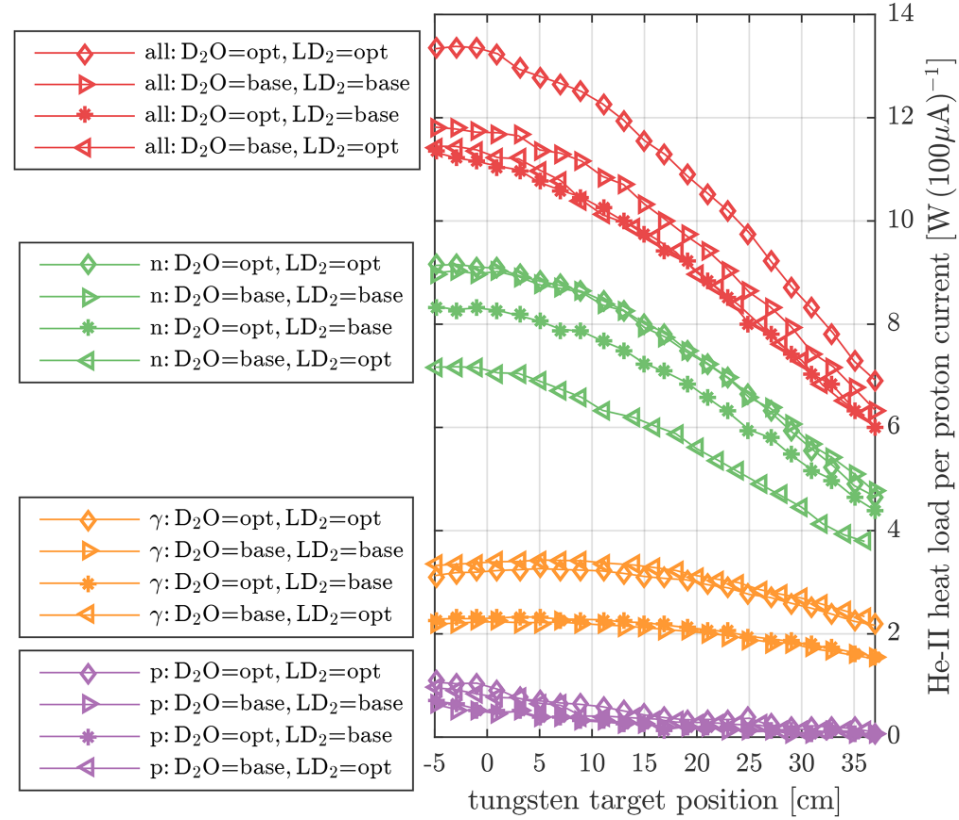


## Cold neutron spectrum in He-II volume per 80 kW protons



UCN production better optimized per He-II heat load

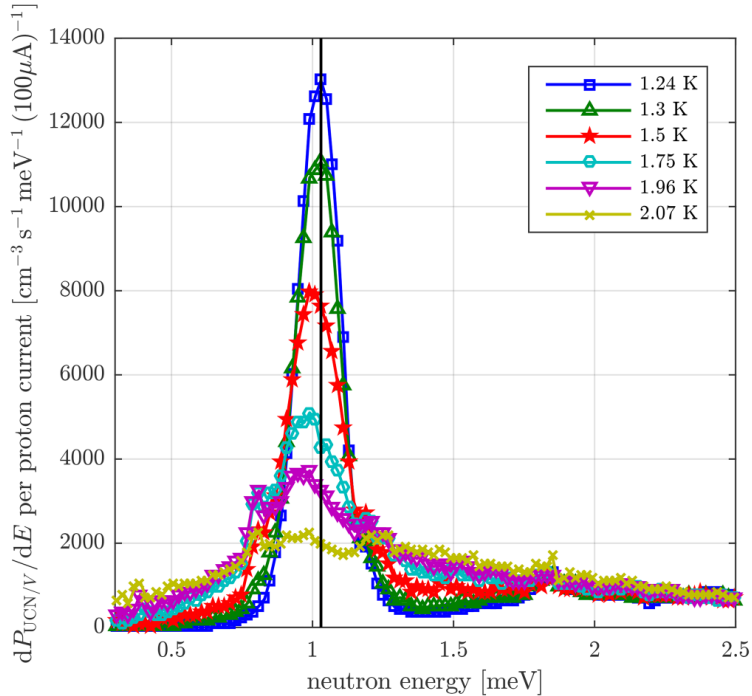
## He-II heat load per 80 kW protons:



The neutron heating dominated by > 80 meV (epithermal to fast)

\*legend are the D2O premoderator and LD2 moderator thicknesses at their "base" and "optimized" values

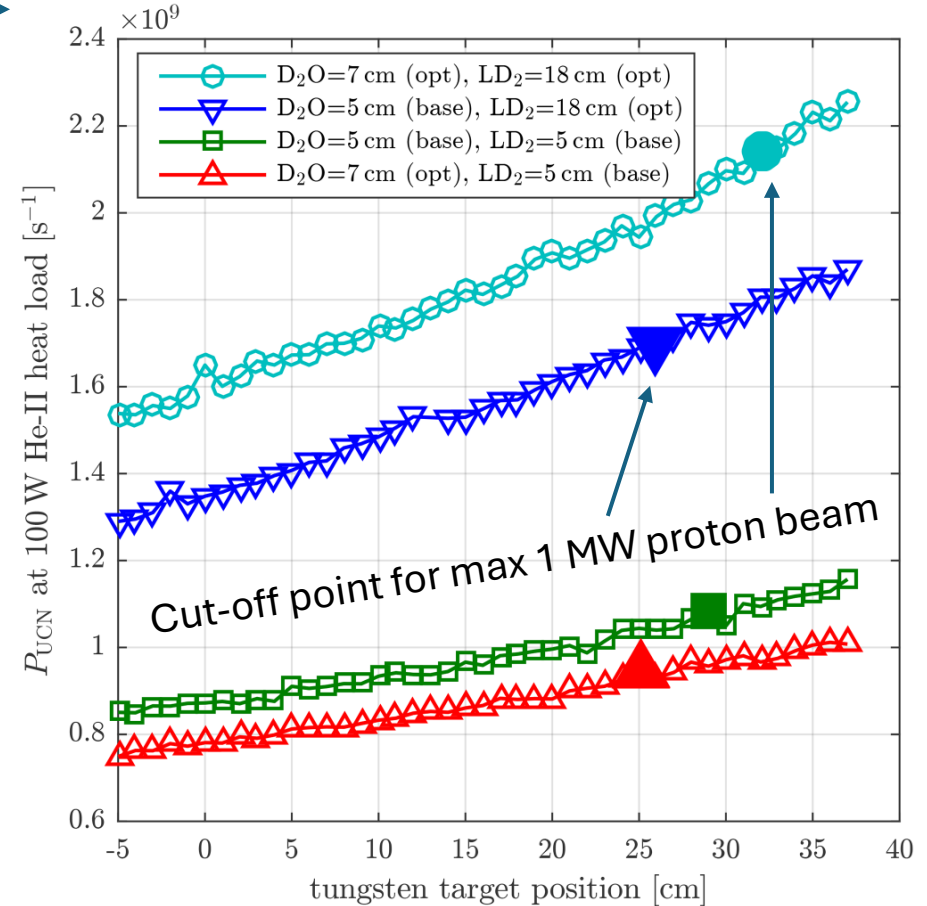
## The CN-energy-differential UCN production rate (i.e. with multi-phonon UCN production)



$^{58}\text{Ni}$  (335 neV)  
minus SF4He  
(18.5 neV)  $E_{\text{ucn}}$   
cut-off energy



## The total UCN production rate



Thanks to Ken Andersen's data from 1994.

- The total **UCN production rate** for **100 W** of heat in the **SF-4He** above is **2.1E9 UCN/s**.
- Include some additional losses\* of 15% and get **1.8 x 10<sup>9</sup> UCN/s**

(\*1 bar pressurized He-II to suppress bubble formation + lack of He-II kernel in MCNP)

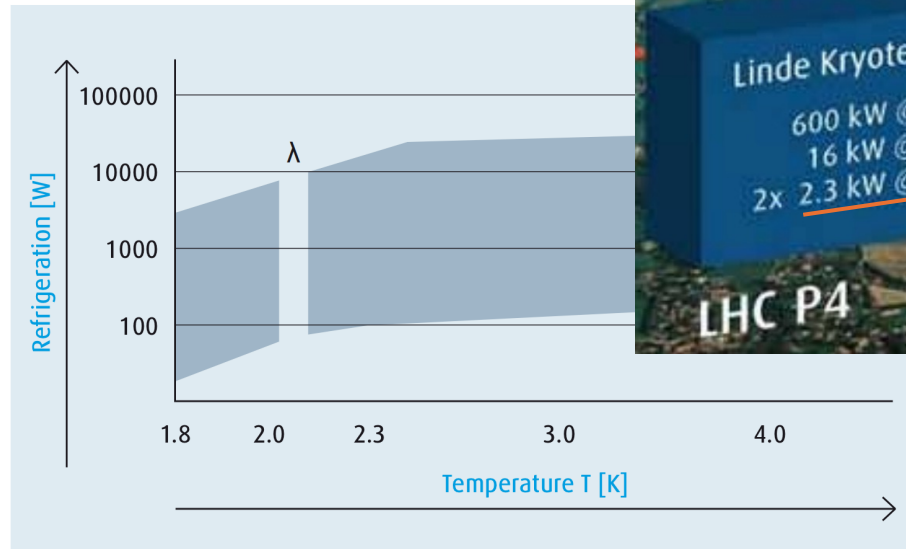
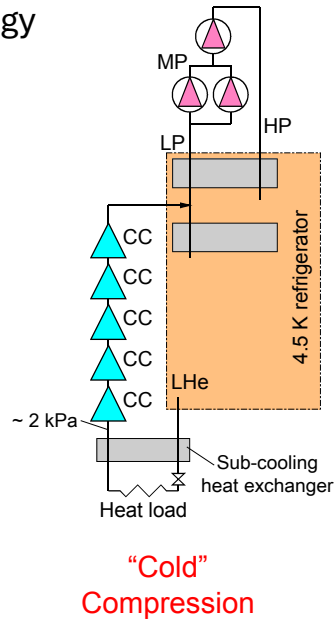
# Sub-cooled $^4\text{He}$ technology



## Superfluid helium refrigeration system

Used at Jefferson Lab, CERN, Fermilab, ESS, etc. for cryogenic superconducting cavities & magnets

Key enabling technology are cold compressors

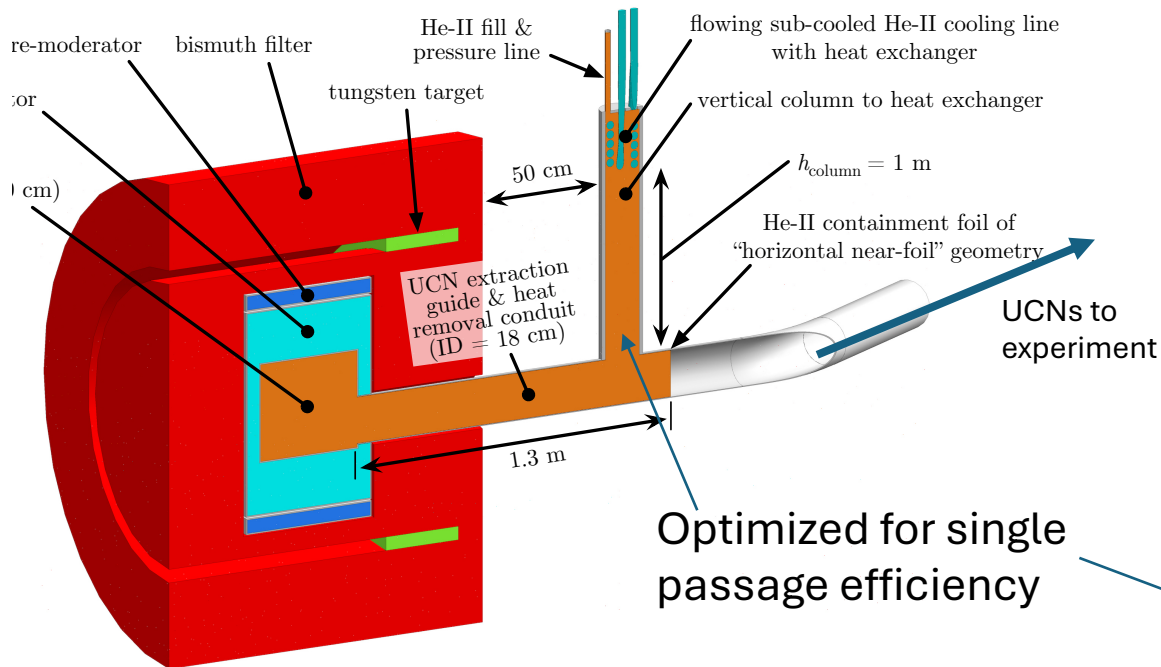


**2.3 kW @ 1.8 K**  
**deployed 20 years**  
**ago. (SNS has a 2.1 K**  
**system)**

- Generally,  $\sim 1.8$  K lowest temp. Requires 4 cold compressors in series. Up to 3 kW of cooling!
- To reach to 1.6K need to go from 4 to 5 cold compression stages, which requires some R&D, which can be off-loaded to external company.
- Note: CERN's Future Circular Collide also requires 1.6 K with kW cooling (we require only 100 W!).

# Single-passage UCN extraction efficiency

- 18 cm  $\varnothing$  UCN extraction guide allows 100 W heat extraction to produce  $\Delta T < 50$  mK. (Our heat flux falls in the Gorter-Mellink regime, mutual friction between normal and superfluid components)
- Tee to heat exchanger. UCNs that reach heat exchanger assumed to be 100% loss (pessimistic)
- Add non-specular reflections in some places to extract UCNs faster
- He-II contained with polypropylene foil supported by a grid (transmission through foil  $\sim 68\%$ ).
- f-factor = W/U of UCN guides assumed to be  $5E-4$ . 3% Lambertian diffuse. The up-scattering in He-II loss is  $\sim 55\%$  of total.



**TABLE III.** Summary of steps taken to reach  $\epsilon_{\text{tot single}} = \epsilon_{\text{sim}} \epsilon_{\text{grid}} \epsilon_{\text{guide}} = 26\%$  (shown in bold) for the horizontal near-foil UCN extraction geometry for  $T = 1.6$  K.

Configuration	$\epsilon$
Baseline (ideal Al foil, $P_{\text{diffuse}} = 3\%$ everywhere)	35% ( $\epsilon_{\text{sim}}$ )
Add diffuse reflections in converter volume ( $P_{\text{diffuse}} = 50\%$ )	43% ( $\epsilon_{\text{sim}}$ )
Add diffuse reflections in vertical column ( $P_{\text{diffuse}} = 50\%$ )	45% ( $\epsilon_{\text{sim}}$ )
Switch from ideal Al foil (54 neV) to ideal PP ( $-8$ neV)	53% ( $\epsilon_{\text{sim}}$ )
Add more realistic PP elastic scattering ( $\lambda_{\text{scat}} = 20 \mu\text{m}$ )	36% ( $\epsilon_{\text{sim}}$ )
Include PP foil support grid loss ( $\epsilon_{\text{grid}} = 90\%$ )	32% ( $\epsilon_{\text{sim}} \epsilon_{\text{grid}}$ )
Include 4 m guide loss to external volume ( $\epsilon_{\text{guide}} = 80\%$ )	<b>26% (<math>\epsilon_{\text{tot single}}</math>)</b>

# UCN current out of source


- For our  **$1.8 \times 10^9$  UCN/s total production rate**. At the end of an 18 cm diameter guide 4 m away from the source (e.g. outside biological shielding) the UCN current becomes  **$5 \times 10^8$  UCN/s**.

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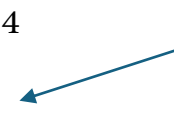
$V_{\text{bottle}}$ (l)	5	50	500	$5 \times 10^3$	$5 \times 10^4$
$\rho_{\text{bottle}}$ ( $\times 10^4$ UCN $\text{cm}^{-3}$ )	1.12	1.11	1.05	0.80	0.31
$\tau_{\text{bottle}}$ (s)	0.11	1.1	10	80	315

UCN density loaded into external experimental volume



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- High-current UCN sources are also ideal for producing a high current of **Very Cold Neutrons** (VCNs). (Need more simulations, rough estimate  $\sim 10^2$  higher than current strongest VCN beam)

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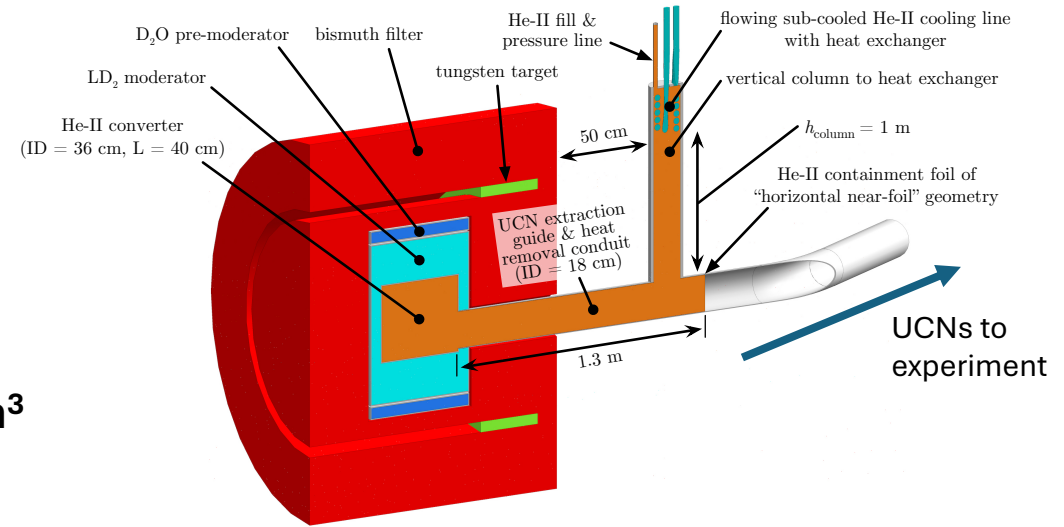
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- Single-passage optimized sources (and assuming no return during filling) are less sensitivity to variations in UCN guide losses, especially difficult for cryogenic guides. Depending on the geometry, when using a density-optimized source to fill an external volume UCNs have to make several passages from source to volume. The transport extraction efficiency becomes  $\sim (\epsilon_{\text{single passage}})^{\text{average no. of passages}}$

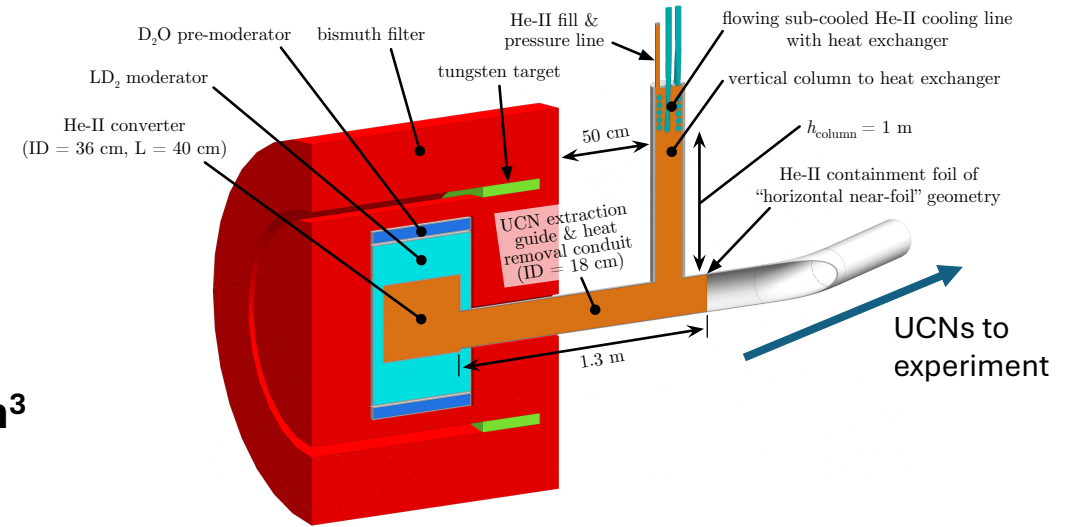
# Summary

- 40 L vessel He-II @ 1.6 K with 100 W cooling
- Can handle 1 MW proton beam power
- $1.8 \times 10^9$  UCN/s production rate
- Optimize for single-passage extraction to get  **$5 \times 10^8$  UCN/s current 4 m away**
- Can deliver to experiments  **$\sim 10,000$  UCN/cm<sup>3</sup>** (or 3,000 UCN/cm<sup>3</sup> for 160 neV max UCN energy)
- Big advantage: **filling large volumes or many experiments simultaneously** (or a flow-through UCN experiments)



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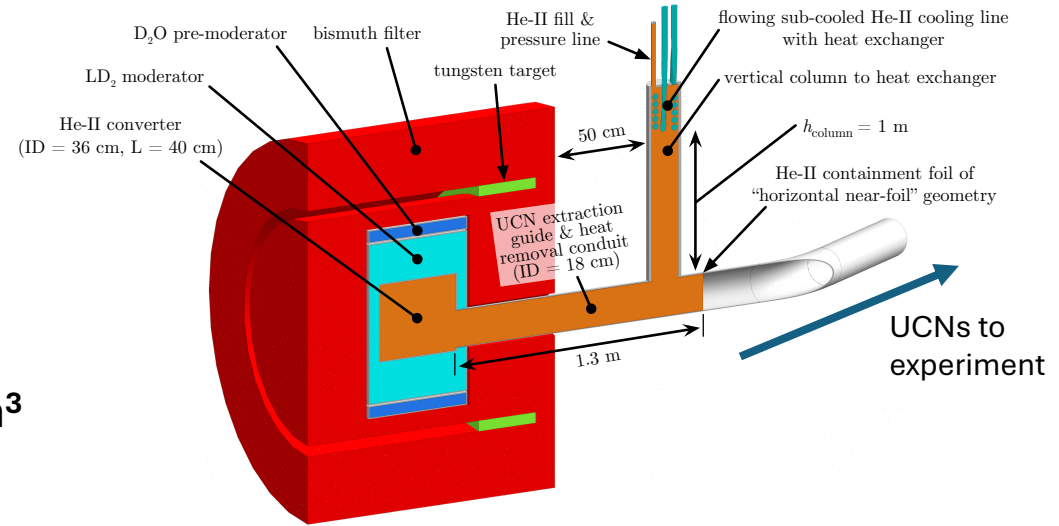
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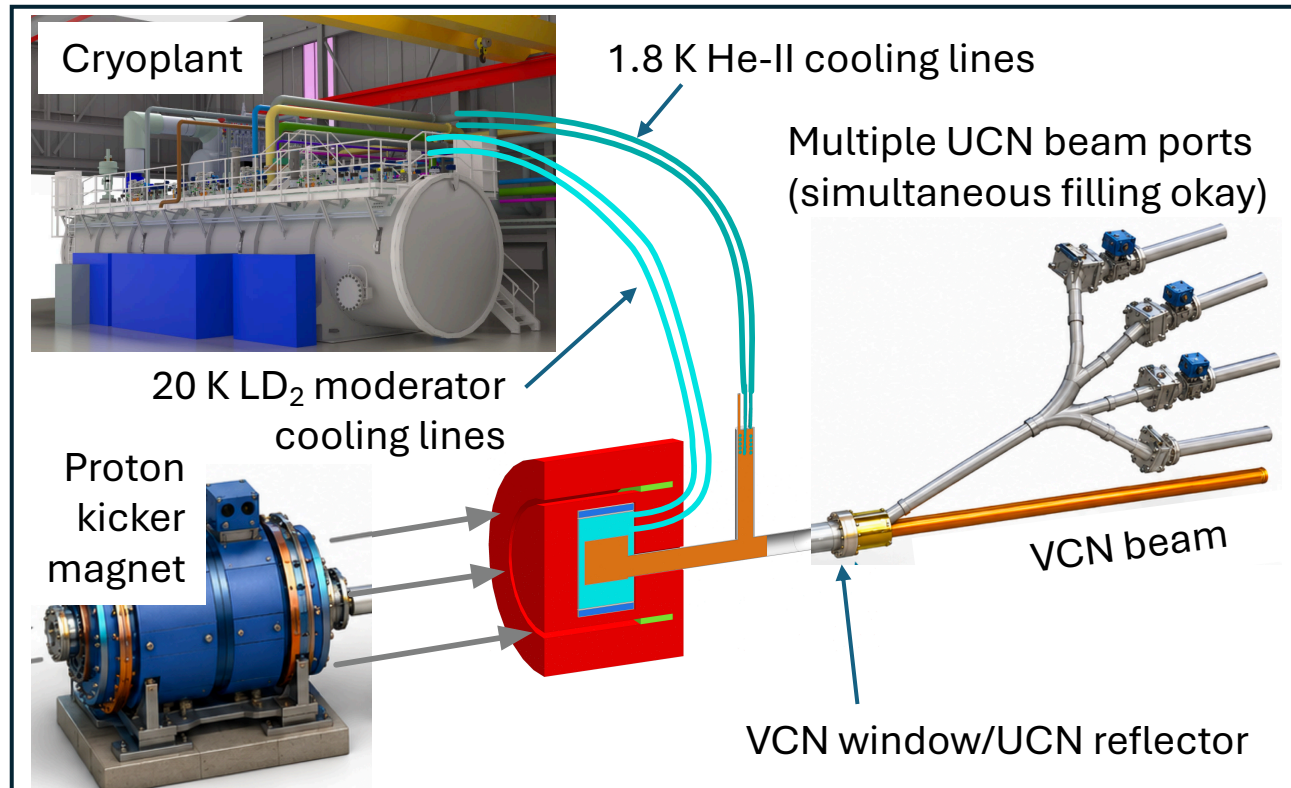
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- nEDM, neutron lifetime,  $\beta$ -decay, quantum levels experiments would all reach a new level. Many of these are “shovel-ready” and can be brought here and run immediately. New previously not-possible experiments could then emerge.

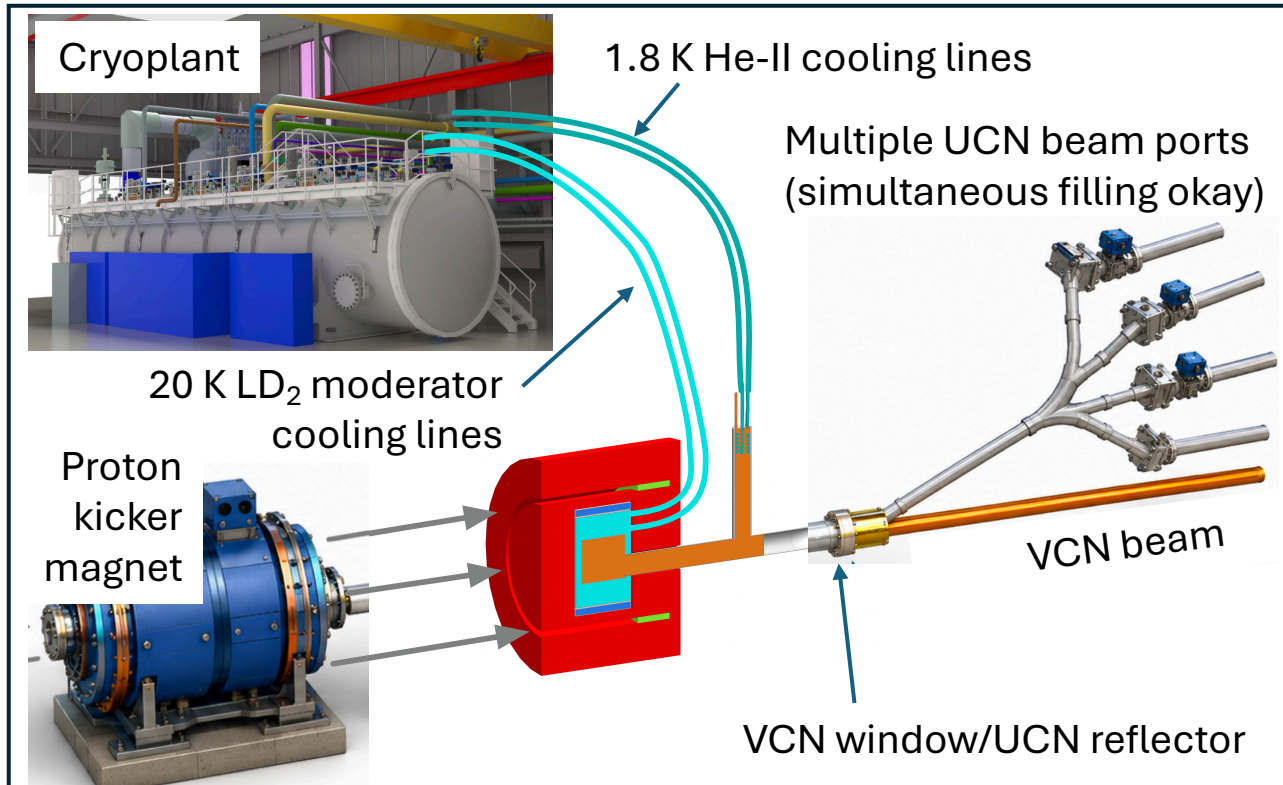
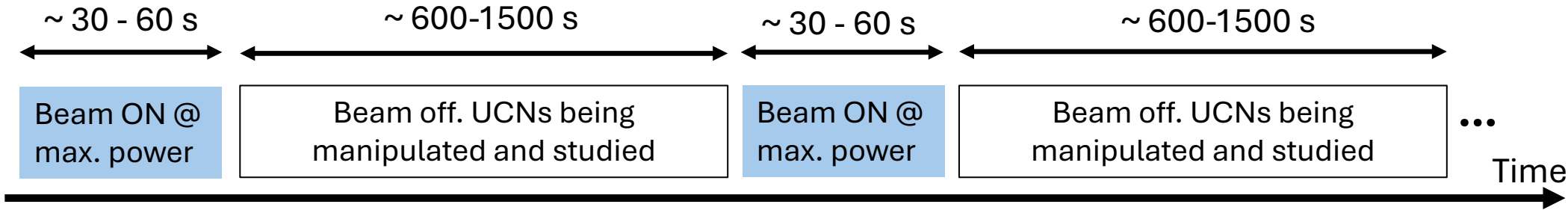


# UCN source facility and proton beam use



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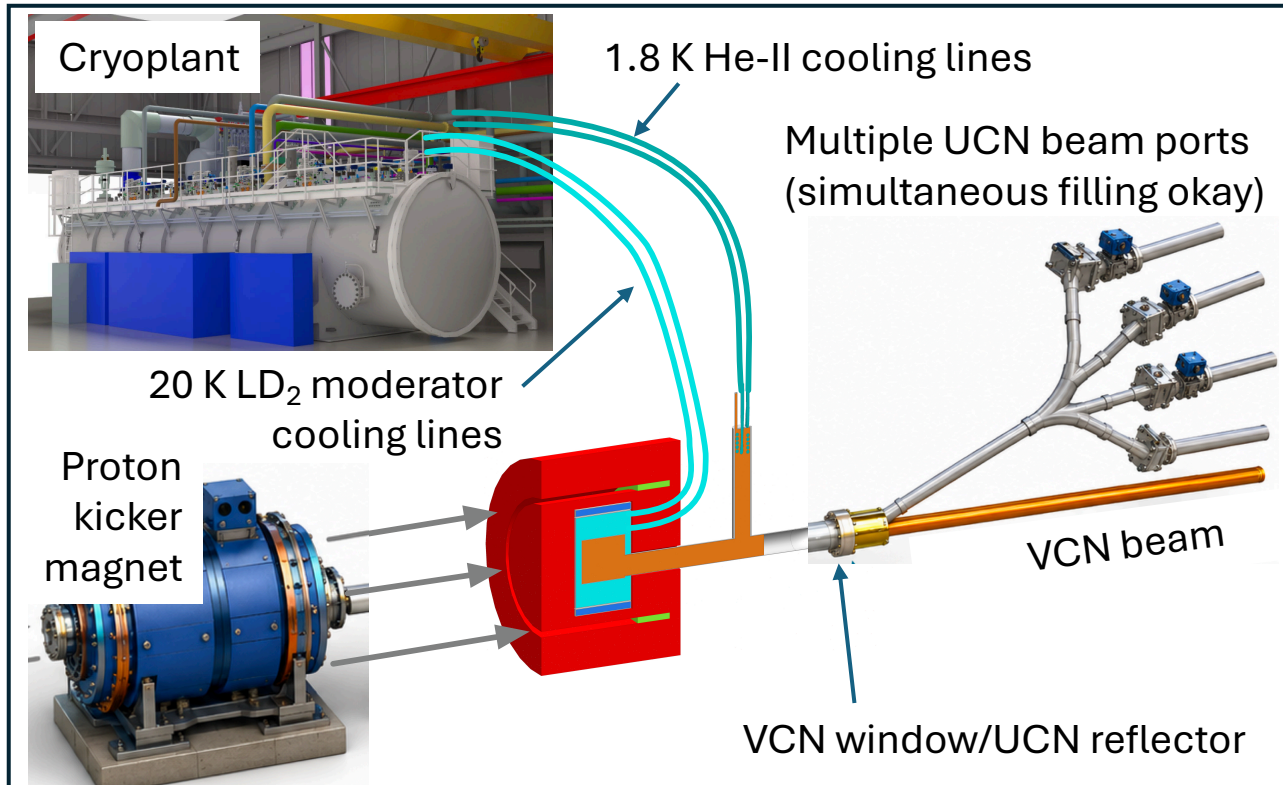
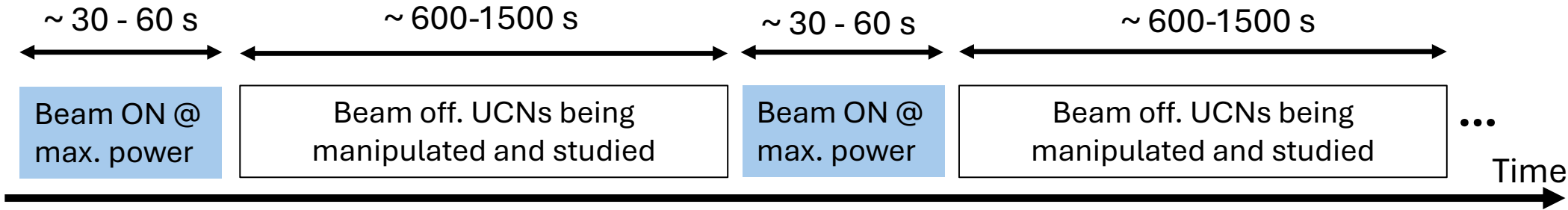
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# UCN source facility and proton beam use

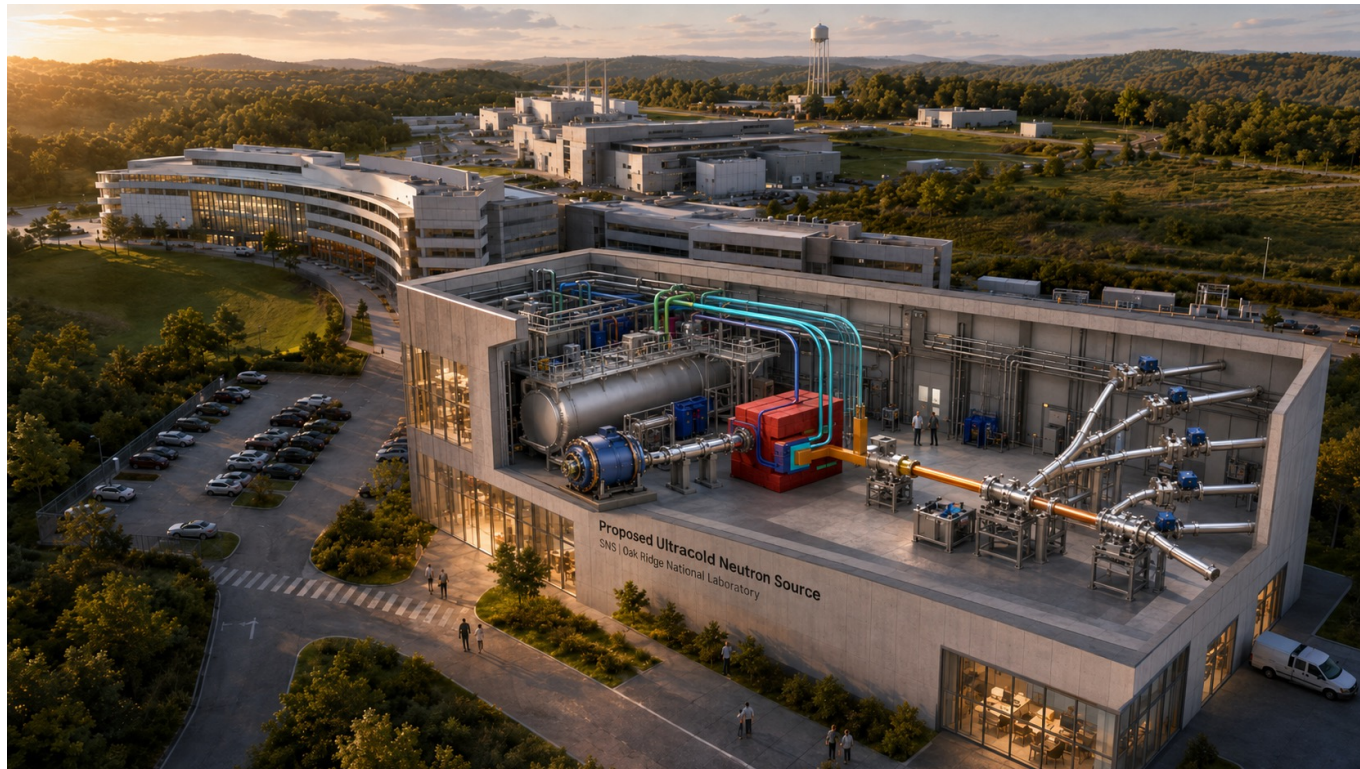
## Proton beam use time structure



- This “fill and empty/study” time structure **used in most UCN experiments**
- Density source requires very **short fill times**
- This UCN facility would only use **2-10% time-averaged power.**
- **Great for sharing with other facilities!**

# Thank you!

ChatGPT thinks this UCN source facility would look something like this at the SNS....



Extra slides

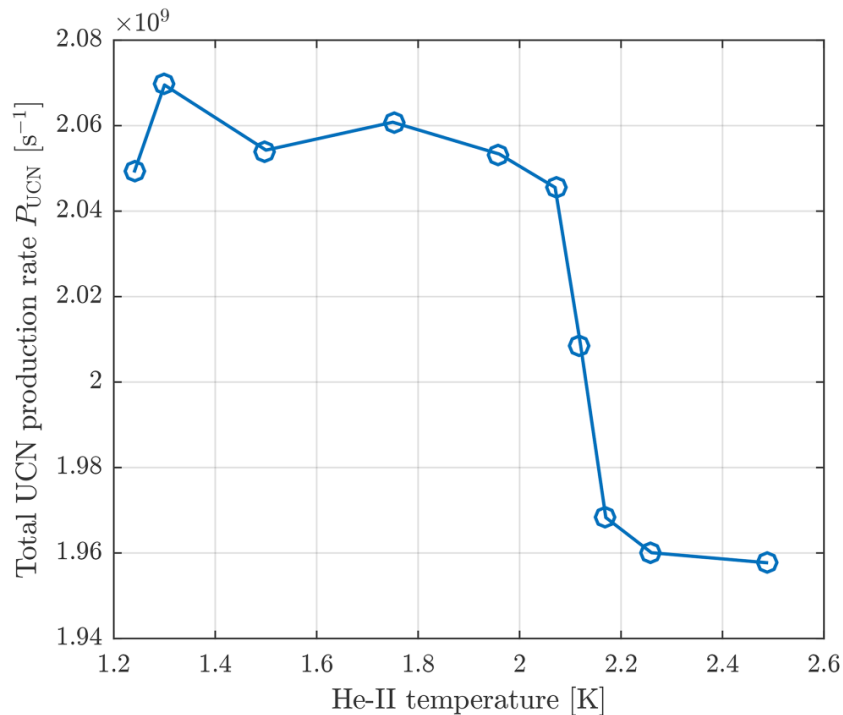
# What if we use warmer SF-<sup>4</sup>He in a UCN source?

PHYSICAL REVIEW C **93**, 025501 (2016)

## Ultracold-neutron production and up-scattering in superfluid helium between 1.1 K and 2.4 K

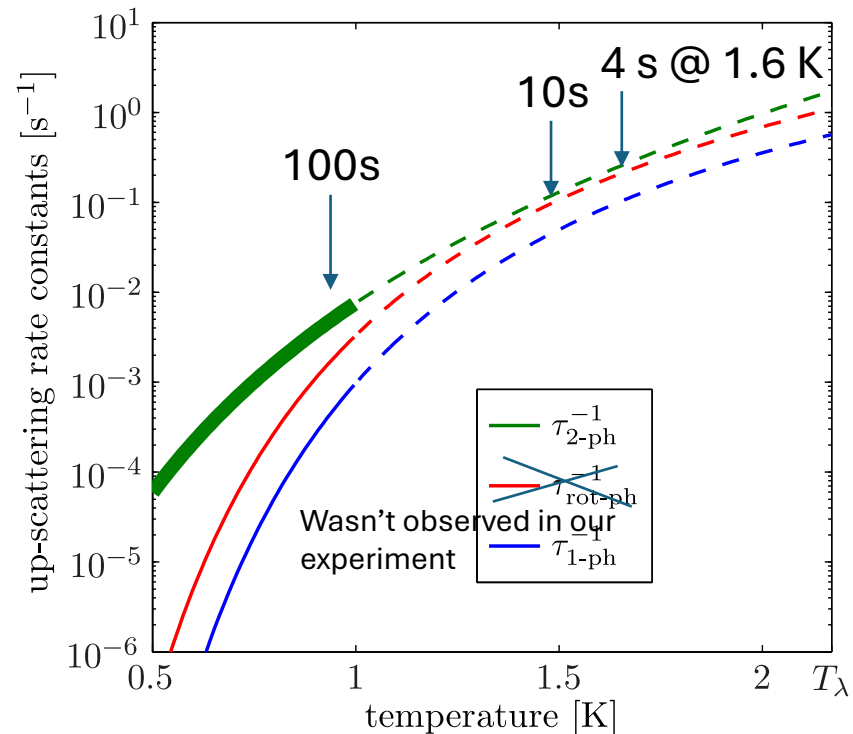
K. K. H. Leung,<sup>1,2</sup> S. Ivanov,<sup>1</sup> F. M. Piegsa,<sup>1,3</sup> M. Simson,<sup>1</sup> and O. Zimmer<sup>1</sup>

### Little impact on UCN production rate



(This is for our CN spectrum)

### UCN “up-scattering loss” increases



# Summary and additional effects on total UCN production rate

Geometry, configuration or effect	Proton power at 100 W He-II (kW)	Neutron heating (%)	Photon heating (%)	Proton heating (%)	CN flux at 1 meV per proton [ $\text{cm}^{-2} \text{s}^{-1} \text{meV}^{-1}$ ( $100 \mu\text{A})^{-1}$ ]	Peak CN flux (meV)	$P_{\text{UCN}}$ ( $\text{s}^{-1}$ )
Mark-3-inspired UCN source (Sec. III)	120	67	28	5	$5.8 \times 10^{10}$	2.6	$0.1 \times 10^9$
Baseline Inverse Geometry (Sec. IV)	300				$4.5 \times 10^{10}$	2.2	$0.2 \times 10^9$
<i>Inverse Geometry after modifications: (Secs. V and VI)</i>							
D <sub>2</sub> O = 5 cm (b), LD <sub>2</sub> = 5 cm (b), targ = 0 cm (b)	680	77	19	4	$7.4 \times 10^{10}$	2.0	$0.7 \times 10^9$
D <sub>2</sub> O = 5 cm (b), LD <sub>2</sub> = 18 cm (o), targ = 0 cm (b)	700	63	30	7	$11 \times 10^{10}$	2.0	$1.3 \times 10^9$
D <sub>2</sub> O = 7 cm (o), LD <sub>2</sub> = 5 cm (b), targ = 0 cm (b)	710	74	21	5	$6.0 \times 10^{10}$	2.0	$0.8 \times 10^9$
D <sub>2</sub> O = 7 cm (o), LD <sub>2</sub> = 18 cm (o), targ = 0 cm (b)	600	69	24	7	$17 \times 10^{10}$	1.7	$1.6 \times 10^9$
D <sub>2</sub> O = 5 cm (b), LD <sub>2</sub> = 5 cm (b), targ = 29 cm (o)*	1000*	76	23	1	$6.3 \times 10^{10}$	2.0	$1.1 \times 10^9$
D <sub>2</sub> O = 7 cm (o), LD <sub>2</sub> = 5 cm (b), targ = 25 cm (o)*	1000*	74	24	2	$5.3 \times 10^{10}$	2.1	$0.9 \times 10^9$
D <sub>2</sub> O = 5 cm (b), LD <sub>2</sub> = 18 cm (o), targ = 26 cm (o)*	1000*	62	36	2	$9.9 \times 10^{10}$	2.0	$1.7 \times 10^9$
D <sub>2</sub> O = 7 cm (o), LD <sub>2</sub> = 18 cm (o), targ = 32 cm (o)*	1000*	67	32	1	$14 \times 10^{10}$	1.7	$2.1 \times 10^9$
MCNP He-II kernel (10% reduction, Sec. VI D)							<u><math>1.9 \times 10^9</math></u>
He-II pressure at 1 bar (3% reduction, Sec. VI E)							<u><math>1.8 \times 10^9</math></u>

- MCNP did not have a superfluid 4He scattering kernel. Since our paper, we have developed this kernel with Chris Lavelle & Takeyasu Ito
- We want to pressurize SF4He to a modest 1 bar to reduce bubble formation that can scatter UCNs